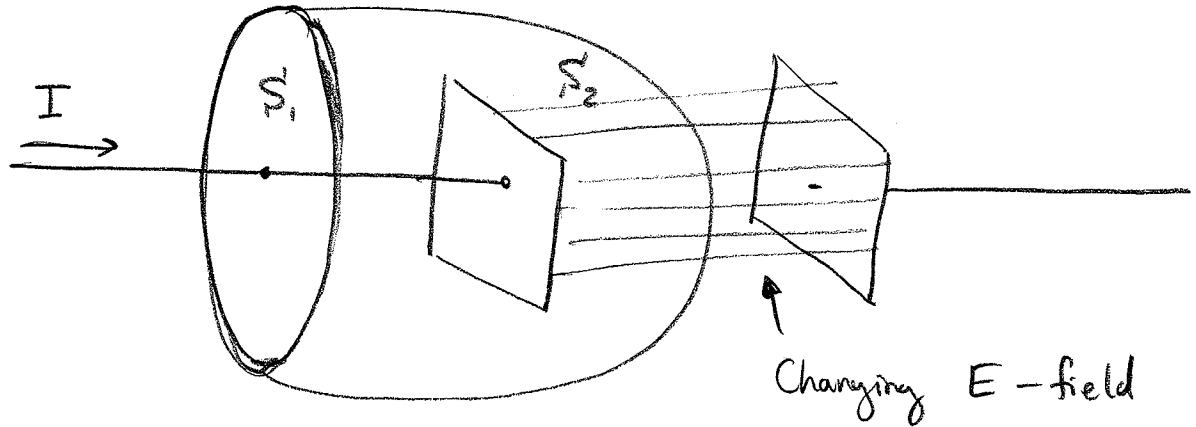


# DISPLACEMENT CURRENT & MAXWELL'S EQS

(114)

Ampere's Law may fail when applied to non constant currents



$$\partial S_1 = \partial S_2 = C$$

$$\int_C \vec{B} \cdot d\vec{l} = \mu_0 I_{S_1} \neq \mu_0 I_{S_2}$$

current through  $S_1$  is  $I$

zero current through  $S_2$

HOWEVER there is a changing E-field through this  $S_2$ .

Thus, Maxwell suggested that the correct form of Ampere's Law is

$$\int_{\partial S} \vec{B} \cdot d\vec{l} = \mu_0 I_S + \mu_0 I_D$$

both through S

$$I_S = \int_S \vec{J} \cdot d\vec{A}$$

$$I_D = \int_S \epsilon_0 \frac{\partial \vec{E}}{\partial t} \cdot d\vec{A}$$

$$\text{Displacement "Current"} = \epsilon_0 \frac{\partial}{\partial t} \int_S \vec{E} \cdot d\vec{A} = \epsilon_0 \frac{\partial \Phi_E}{\partial t}$$

Remark: for the capacitor counter-example to naive Ampere's law one can show  $I_s = I_0|_{s_2}$   
See Tipler Example 30.1 pg. 1032

The differential form of Ampere's Law with Maxwell's Curvature

$$\int_{\partial S} \vec{B} \cdot d\vec{l} = \mu_0 I_s + \mu_0 I_D$$

$$= \mu_0 \left( \int_S \vec{J} \cdot d\vec{A} + \epsilon_0 \frac{\partial}{\partial t} \int_S \vec{E} \cdot d\vec{A} \right)$$

$$= \int_S \left[ \mu_0 \vec{J} + \epsilon_0 \mu_0 \frac{\partial \vec{E}}{\partial t} \right] \cdot d\vec{A}$$

However,  $\int_{\partial S} \vec{B} \cdot d\vec{l} = \int_S (\nabla \times \vec{B}) \cdot d\vec{A}$  by Stoke's Th<sup>m</sup>. Hence, as this applies for essentially arbitrary surfaces  $S'$ . Finally we find Maxwell's Eq<sup>s</sup>

$$\nabla \times \vec{B} = \mu_0 \vec{J} + \epsilon_0 \mu_0 \frac{\partial \vec{E}}{\partial t}$$

$$\nabla \cdot \vec{E} = \rho / \epsilon_0$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Electromagnetism expressed as a local theory

$$\vec{J} = \frac{\text{current}}{\text{area}}, \quad \rho = \frac{\text{charge}}{\text{volume}}$$

$$\vec{J} = \frac{d\vec{I}}{dA}, \quad \rho = \frac{dQ}{dV}$$

# Electromagnetic Waves in Vacuum

Assume  $\vec{J} = 0$  and  $\rho = 0$  hence

$$\nabla \times \vec{B} = \epsilon_0 \mu_0 \frac{\partial \vec{E}}{\partial t} \qquad \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\nabla \cdot \vec{E} = 0 \qquad \nabla \cdot \vec{B} = 0$$

Consider that,

$$\nabla \times (\nabla \times \vec{E}) = \nabla (\nabla \cdot \vec{E}) - \nabla^2 \vec{E} = -\nabla^2 \vec{E}$$


$$\nabla \times \left(-\frac{\partial \vec{B}}{\partial t}\right) = -\frac{\partial}{\partial t} (\nabla \times \vec{B}) = -\frac{\partial}{\partial t} \left(\epsilon_0 \mu_0 \frac{\partial \vec{E}}{\partial t}\right) = -\epsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\therefore \boxed{\nabla^2 \vec{E} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}} \quad \text{WAVE EQN!}$$

identify as  $\frac{1}{v^2} = \mu_0 \epsilon_0$

$$\therefore v = \frac{1}{\sqrt{\mu_0 \epsilon_0}} = c$$

speed of light in vacuum.

 LIGHT IS AN ELECTROMAGNETIC WAVE!

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# PROPERTIES OF LIGHT

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• Should read your textbook and/or the powerpts. we discussed in lecture. I'll just remind a few essential equations.

• The  $\vec{E}$  &  $\vec{B}$  in light are  $\perp$  in vacuum.

•  $E = cB$  where  $c = \text{speed of light} = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \approx 3 \times 10^8 \frac{\text{m}}{\text{s}}$   
magnitudes (not direction!)

•  $\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B}$  ← Poynting vector, points in direction of motion of the Electromagnetic Wave.

•  $P = \frac{S}{c} = \frac{1}{\mu_0 c} \|\vec{E} \times \vec{B}\| =$  pressure of light

gives rise to radiation pressure from light

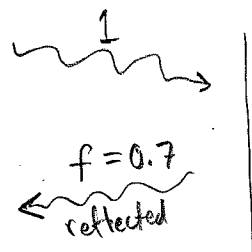
Recall  $P = \frac{\text{FORCE}}{\text{AREA}} = \frac{dE}{dA}$  (generally, but

we'll probably only work constant  $P$  problems)

**E87** (24.4 on p. 821 of Serway)

Consider laser pointer with 3.0 mW beam with area with 2.0 mm = diameter. Find radiation pressure on screen which reflects 70% of light.

$$\text{Intensity} = \frac{\text{Power}}{\text{area}} = \frac{3.0 \text{ mW}}{\pi(1.0 \text{ mm})^2} = 9.6 \times 10^2 \frac{\text{W}}{\text{m}^2} \quad \left( \frac{\text{energy}}{\text{area}(\text{time})} \right)$$



$$P_{\text{avg}} = \frac{S_{\text{avg}}}{c} + \frac{f S_{\text{avg}}}{c} = (1+f) \frac{S_{\text{avg}}}{c} = \frac{(1.7)(9.6 \times 10^2 \frac{\text{W}}{\text{m}^2})}{3 \times 10^8 \text{ m/s}}$$

$$\Rightarrow P = 5.4 \times 10^{-6} \frac{\text{N}}{\text{m}^2}$$

Remark: The magnitude of  $\vec{S}$  is  $S = I = \text{intensity}$  (118)  
 this follows from dimensional analysis among other things.

$$\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B} \sim \frac{c}{\mu_0} B^2$$

$$\Rightarrow \frac{1}{c} S \sim \frac{B^2}{2\mu_0} \sim \frac{\text{energy}}{\text{volume}}$$

$$\text{So } \frac{B^2}{2\mu_0} \cdot c \sim S \sim \frac{\text{energy}}{\text{volume}} \frac{\text{distance}}{\text{time}} \sim \frac{\text{energy}}{\text{time}} \frac{\text{distance}}{\text{volume}}$$

$\therefore$  units for  $S$  are  $\frac{\text{Power}}{\text{area}}$  which is intensity

This is one reason to include  $\frac{1}{\mu_0}$  in def<sup>n</sup> of  $\vec{S}$ . It forces  $|\vec{S}| = \text{intensity}$  along  $\hat{S}$ .

### Isotropic Radiation

spreads out in same way for all directions.  
 If we assume conservation of energy in the wave then we can derive the inverse square law.

### Malus' Law

$$I = I_0 \cos^2 \theta$$

exit intensity

initial intensity

angle of linearly polarized light to the polarizer.

(see pg. 825 of Serway)