

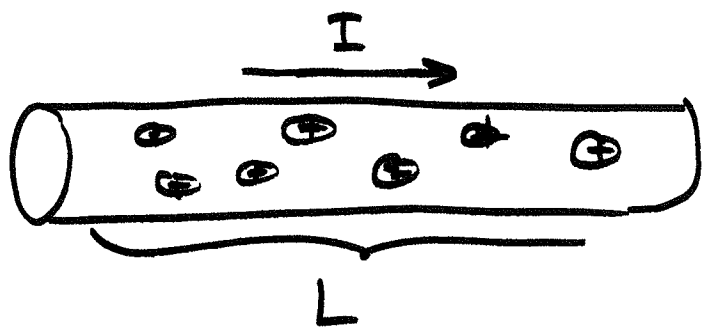
Lecture Notes

$$\vec{F} = q \vec{v} \times \vec{B}$$

Units for $|\vec{B}|$ are Teslas

$$1 T = \frac{N}{C \cdot m/s}$$

1 Gauss ($G = 10^{-4} T$)



$$\vec{I} = q_{total} \vec{v}_{drift}$$

↑
proportional to l

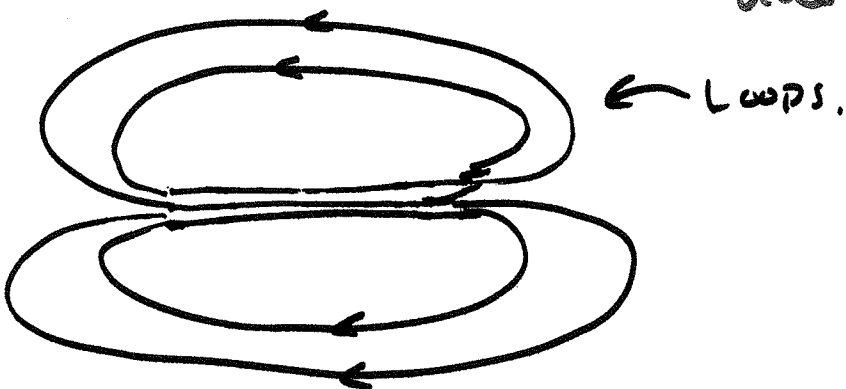
$$\vec{F} = I(\vec{l} \times \vec{B})$$

$$= l(\vec{I} \times \vec{B}) \leftarrow \text{force on the length } l \text{ of wire.}$$

Infinitesimally,

$$d\vec{F} = I d\vec{l} \times \vec{B}$$

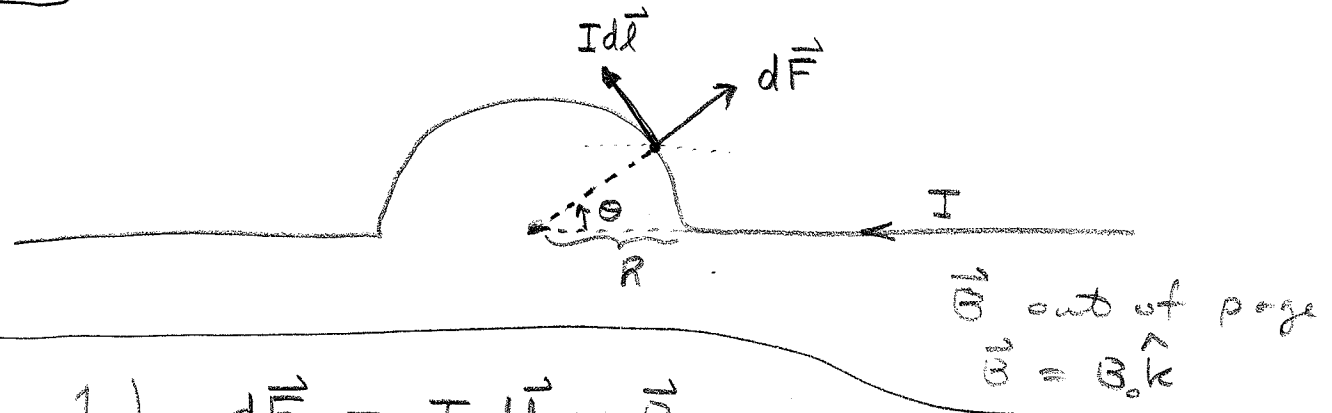
force on $d\vec{l}$
due to \vec{B}



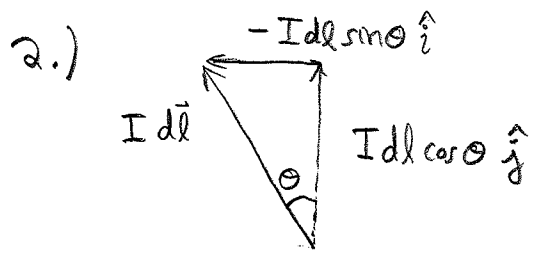
← Loops.

ES9

Find Force on half-loop due to \vec{B}



1.) $d\vec{F} = I d\vec{l} \times \vec{B}$



$$d\vec{F} = (-I \sin\theta dl \hat{i} + I \cos\theta dl \hat{j}) \times (B_0 \hat{k})$$

3.) $\hat{i} \times \hat{k} = -\hat{j}$, $\hat{j} \times \hat{k} = \hat{i}$ and $dl = R d\theta$

$$d\vec{F} = -IR \sin\theta B_0 d\theta (\hat{i} \times \hat{k}) + IR \cos\theta d\theta (\hat{j} \times \hat{k}) B_0$$

$$= (IR \sin\theta B_0 d\theta) \hat{j} + (IR \cos\theta B_0 d\theta) \hat{i}$$

4.) integrate for $0 \leq \theta \leq \pi$

$$\vec{F} = \left(IR B_0 \int_0^\pi \sin\theta d\theta \right) \hat{j} + \left(IR B_0 \int_0^\pi \cos\theta d\theta \right) \hat{i}$$

$$= \boxed{2IR B_0 \hat{j}}$$

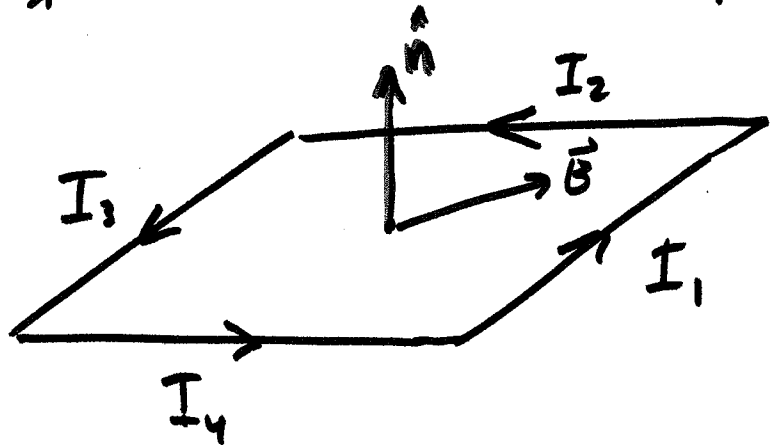
• velocity selector

$$\vec{F} = q\vec{E} + q\vec{v} \times \vec{B}$$

need $v = \frac{E}{B}$ to balance

If $\vec{E} \perp \vec{B}$

• torque on current loop. (for axis through center of loop in plane of current)



$$\vec{\tau} = \vec{\mu} \times \vec{B}$$

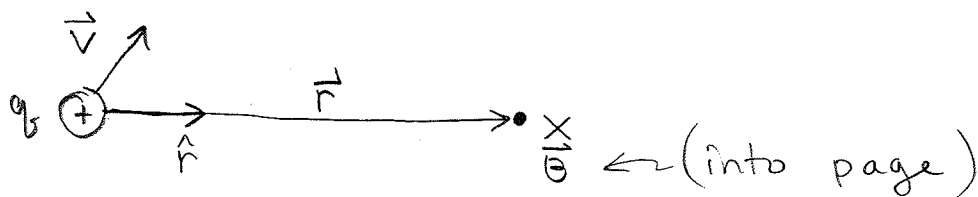
where $\vec{\mu} = NIA\hat{n}$

\uparrow # of turns
 \uparrow current
 \uparrow area of loop
 \uparrow normal (ccw) choice.

Magnetic Field due to moving point charge

(85)

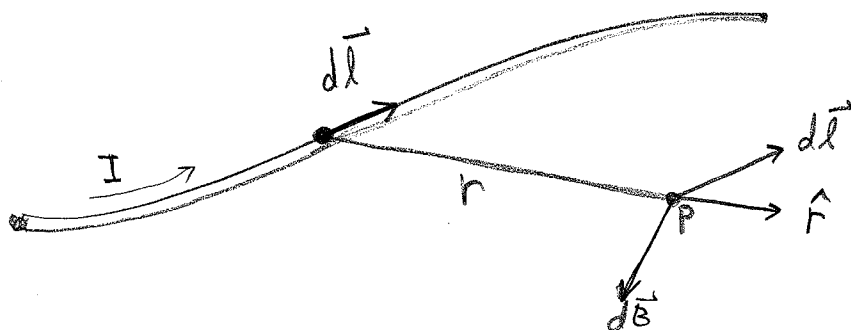
$$\vec{B} = \frac{\mu_0}{4\pi} \frac{q \vec{v} \times \hat{r}}{r^2}$$



$$\mu_0 = 4\pi \times 10^{-7} \frac{\text{T}\cdot\text{m}}{\text{A}} = 4\pi \times 10^{-7} \frac{\text{N}}{\text{A}^2}$$

Biot-Savart Law

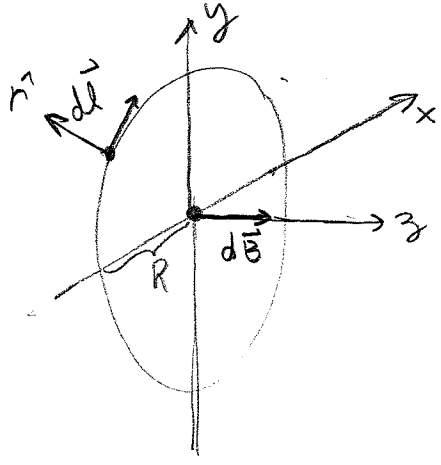
$$d\vec{B} = \frac{\mu_0}{4\pi} \frac{I d\vec{l} \times \hat{r}}{r^2}$$



To find total \vec{B} at P we have to add up all the $d\vec{B}$'s from each infinitesimal current $I d\vec{l}$.

E63

: find \vec{B} due to a loop of current. Find \vec{B} at the center of the loop.



Note: $\hat{r} \times d\vec{l} = dl \hat{k} = R d\theta \hat{k}$

for each pt. around loop always get \hat{k} .

$$\vec{B} = \int_{\text{loop}} d\vec{B} = \int_0^{2\pi} \frac{\mu_0}{4\pi} \frac{I R d\theta \hat{k}}{R^2}$$

$$\frac{I d\vec{l} \times \hat{r}}{r^2} = \frac{R d\theta \hat{k}}{R^2}$$

(in this case)

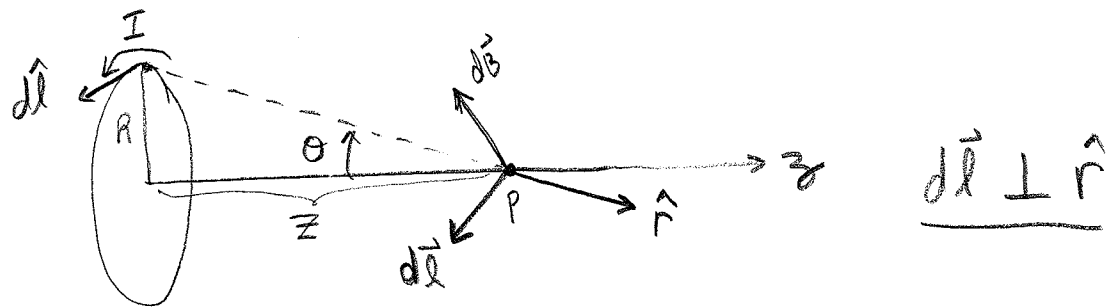
$$\therefore \vec{B} = \frac{\mu_0 I (2\pi) \hat{k}}{4\pi R}$$

$$\Rightarrow \boxed{\vec{B} = \frac{\mu_0 I \hat{k}}{2R}}$$

current at center of loop.

If we had N-loops then $\vec{B} = \frac{\mu_0 N I \hat{k}}{2R}$ by almost the same arguments.

E64 Given loop of current find \vec{B} at p.



Apply Biot-Savart Law,

$$|d\vec{B}| = \frac{\mu_0}{4\pi} \frac{I |d\vec{l} \times \hat{r}|}{r^2} = \frac{\mu_0}{4\pi} \frac{I dl}{(z^2 + R^2)}$$

Symmetry \Rightarrow only $B_z \neq 0$.

$$dB_z = \sin\theta dB = \frac{R}{\sqrt{z^2 + R^2}} \frac{\mu_0 I dl / 4\pi}{z^2 + R^2}$$

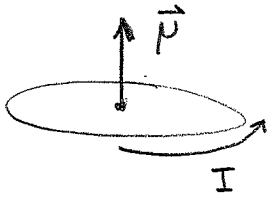
$$B_z = \int_0^{2\pi} \frac{\mu_0 I R^2 d\theta}{4\pi (z^2 + R^2)^{3/2}}$$

$$B_z = \frac{\mu_0 I R^2}{2 (z^2 + R^2)^{3/2}}$$

Remark: returning to E64 we found,



$$|\vec{B}| = B = \frac{\mu_0 I R^2}{2(z^2 + R^2)^{3/2}}$$



If $|z| \gg R$ then $(z^2 + R^2)^{3/2} \rightarrow |z|^3$
 Thus the far-field dipole strength is simply

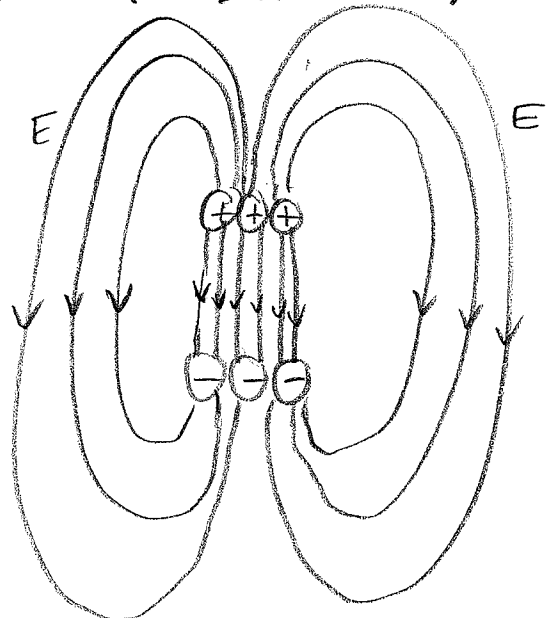
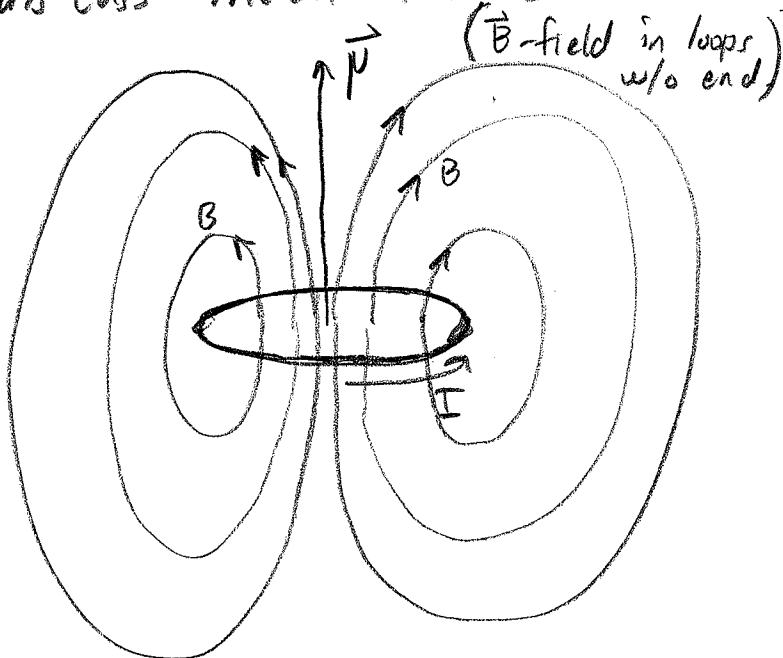
$$B = \frac{\mu_0 I R^2}{2|z|^3}$$

$$= \frac{\mu_0}{4\pi} \frac{2I\pi R^2}{|z|^3}$$

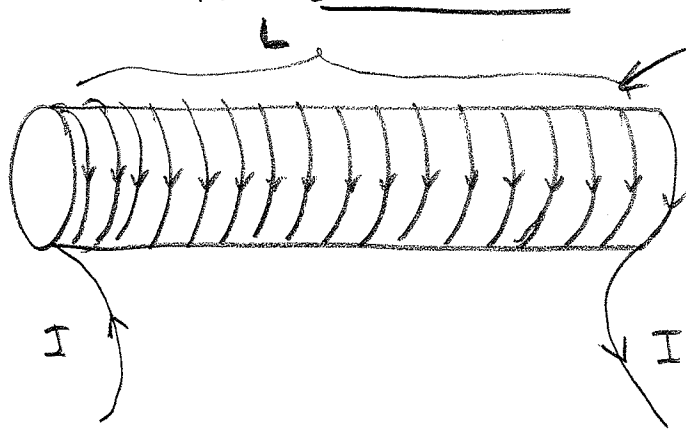
$$= \frac{\mu_0}{4\pi} \frac{2\mu}{|z|^3}$$

Defn/ $\mu = \pi R^2 I$ ← magnitude of dipole moment of current loop

The \vec{E} -field dipole was almost same. $E = \frac{1}{4\pi\epsilon_0} \frac{2dq}{|z|^3}$
 where $dq =$ Electric dipole moment (we didn't discuss much this semester) (\vec{E} -field terminates on source charges)



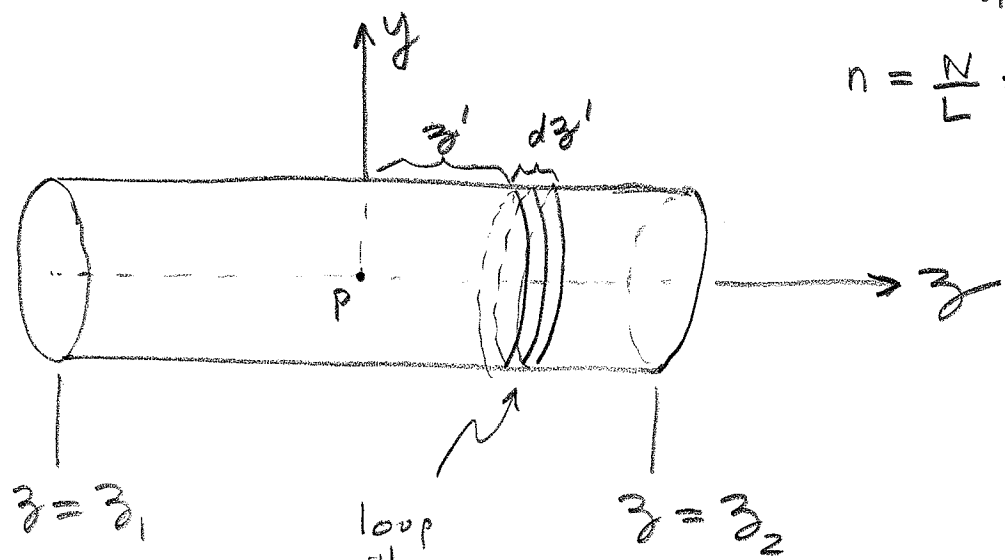
E65 \vec{B} - due to solenoid



tightly wound wire. Current loops glued together intuitively.

$N = \#$ of turns.
for $z_1 \leq z \leq z_2$

$$n = \frac{N}{L} = \frac{\text{turns}}{\text{length}}$$



see E64

$$dz = n I dz' = \left(\frac{\# \text{ turns}}{\text{length}} \right) (\text{current}) (\text{length}) \checkmark$$

$$dB_z = \frac{\mu_0 R^2 n I dz'}{2 [(z - z')^2 + R^2]^{3/2}} \quad \text{checking concept.}$$

$$\therefore B_z = \int_{z_1}^{z_2} \frac{\mu_0 R^2 n I dz'}{2 [(z - z')^2 + R^2]^{3/2}}$$

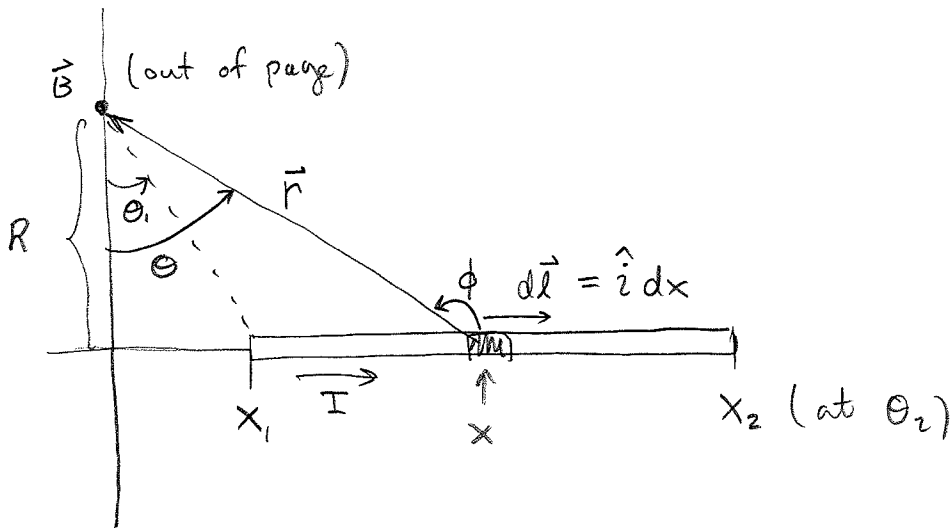
$$B_z = \frac{\mu_0 n I}{2} \left(\frac{z - z_1}{\sqrt{(z - z_1)^2 + R^2}} - \frac{z - z_2}{\sqrt{(z - z_2)^2 + R^2}} \right)$$

Magnetic field on axis of solenoid.

If $L \gg R$ then the solenoid is "long" and

$$B_z \approx \mu_0 n I \quad (\text{center}) \quad B_z = \frac{1}{2} \mu_0 n I \quad (\text{ends})$$

E65 \vec{B} due to segment of wire



$$\begin{aligned}
 dB &= \frac{\mu_0}{4\pi} \frac{I dx}{r^2} \sin\phi \\
 &= \frac{\mu_0}{4\pi} \frac{I dx}{r^2} \cos\theta \\
 &= \frac{\mu_0}{4\pi} \frac{I}{r^2} R \sec^2\theta \cos\theta d\theta \\
 &= \frac{\mu_0 I R}{4\pi r^2} \sec\theta d\theta \\
 &= \frac{\mu_0 I R}{4\pi} \frac{\cos^2\theta}{R^2} \sec\theta d\theta \\
 &= \frac{\mu_0 I}{4\pi R} \cos\theta d\theta
 \end{aligned}$$

Note: $x = R \tan\theta$

$dx = R \sec^2\theta d\theta$

$\therefore \cos\theta = \frac{R}{r}$

$r^2 = R^2 / \cos^2\theta$

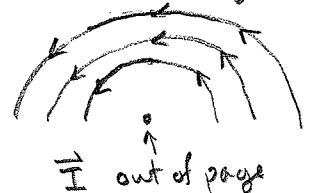
$$B = \int_{\theta_1}^{\theta_2} \frac{\mu_0}{4\pi} \frac{I}{R} \cos\theta d\theta = \frac{\mu_0}{4\pi} \frac{I}{R} (\sin\theta_2 - \sin\theta_1)$$

$$\Rightarrow \underline{B = \frac{\mu_0 I}{2\pi R}}$$

is $\theta_1 \rightarrow 90^\circ$
 $\theta_2 \rightarrow 90^\circ$

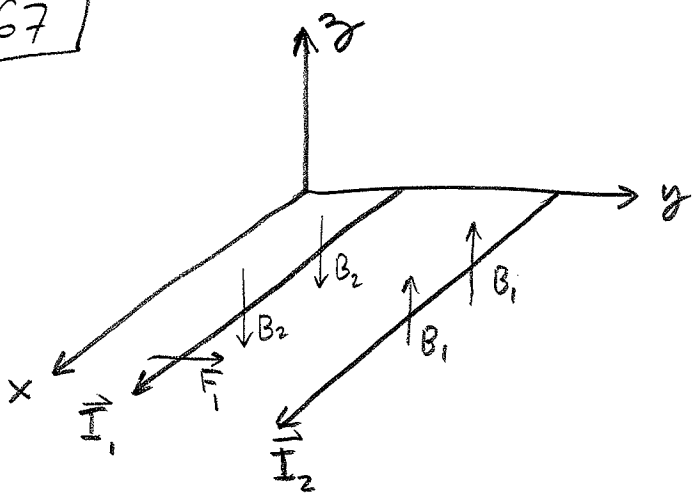
Streamlines of \vec{B} -field

\vec{B} wraps cylindrically around wire as specified by RHR



Force due to parallel currents

E67



\vec{B}_2 is field due to \vec{I}_2
 \vec{B}_1 is field due to \vec{I}_1
 (in my picture)

$$\vec{F}_1 = \vec{F}_{\text{on } I_1 \text{ over } l_1} = l_1 \vec{I}_1 \times \vec{B}_2 \quad (\text{in } +\hat{j} \text{ -direction})$$

$$\vec{F}_2 = \vec{F}_{\text{on } I_2 \text{ over } l_2} = l_2 \vec{I}_2 \times \vec{B}_1 \quad (\text{in } -\hat{j} \text{ -direction})$$

\therefore parallel currents attract.

E68

What if we reverse direction of \vec{I}_2 ?

E69

the force per unit length is found from

$$|d\vec{F}_2| = |I_2 d\vec{l}_2 \times \vec{B}_1| = I_2 dl_2 B_1$$

$$\frac{dF_2}{dl_2} = I_2 B_1 \approx \frac{\mu_0 I_1 I_2}{2\pi R}$$

long wires

Remark: Ampere defined by E67, E68, E69. This in turn defines the Coulomb