

LECTURE 8: MAXWELL'S EQUATIONS & ENERGY / MOMENTUM

(§1.8 & 1.9 Carroll)

①

Maxwell added a term to the Eq's governing electric and magnetic fields which ultimately implied understanding light as an EM wave. It seems Carroll is using units where μ_0 and ϵ_0 are omitted permeability of the vacuum \leftarrow permittivity of the vacuum \leftarrow

$$c = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

$$\nabla \times \vec{B} - \partial_t \vec{E} = \vec{J}$$

$$\tilde{\epsilon}^{ijk} \partial_j B_k - \partial_0 E^i = J^i$$

$$F^{0i} = E^i = -F^{i0}$$

$$\nabla \cdot \vec{E} = \rho$$

$$\partial_i E^i = J^0$$

$$F^{ij} = \tilde{\epsilon}^{ijk} B_k$$

$$\nabla \times \vec{E} + \partial_t \vec{B} = 0$$

$$\tilde{\epsilon}^{ijk} \partial_j E_k + \partial_0 B^i = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\partial_i B^i = 0$$

Define the field strength tensor $F^{\mu\nu}$ and

$$J^\mu = (\rho, J^1, J^2, J^3)$$

is the current

4-vector

\vec{E} = electric field
 \vec{B} = magnetic field

\vec{J} = current density

$\rho = \frac{dq}{dV}$ charge density

$$\begin{aligned} G^{0i} &= \theta^i = -G^{i0} \\ G^{ij} &= \tilde{\epsilon}^{ijk} E_k \\ \partial_\mu G^{\nu\mu} &= 0 \end{aligned}$$

$$\begin{aligned} \partial_j F^{ij} - \partial_0 F^{0i} &= J^i \\ \partial_i F^{0i} &= J^0 \end{aligned}$$

(Ampere's Law with Correction)
 (Gauss' Law)

$$\partial_\mu F^{\nu\mu} = J^\nu$$

then see \rightarrow for the other two Maxwell Eq's

(not in Carroll, but the Dual Tensor $G^{\mu\nu}$ is natural to consider)

$$[F^{\mu\nu}] = \begin{bmatrix} 0 & E^1 & E^2 & E^3 \\ -E^1 & 0 & B^3 & -B^2 \\ -E^2 & -B^3 & 0 & B^1 \\ -E^3 & B^2 & -B^1 & 0 \end{bmatrix}$$

$$\partial_\mu F_{\nu\lambda} = 0$$

$$\partial_\mu F_{\nu\lambda} + \partial_\nu F_{\lambda\mu} + \partial_\lambda F_{\mu\nu} = 0$$

$$[F_{\mu\nu}] = [\eta_{\mu\alpha} \eta_{\nu\beta} F^{\alpha\beta}]$$

$$= \begin{bmatrix} -1 & & & \\ & 1 & & \\ & & 1 & \\ & & & 1 \end{bmatrix} \begin{bmatrix} 0 & E^1 & E^2 & E^3 \\ -E^1 & 0 & B^3 & -B^2 \\ -E^2 & -B^3 & 0 & B^1 \\ -E^3 & B^2 & -B^1 & 0 \end{bmatrix} \begin{bmatrix} -1 & & & \\ & 1 & & \\ & & 1 & \\ & & & 1 \end{bmatrix}$$

$$= \begin{bmatrix} 0 & -E^1 & -E^2 & -E^3 \\ -E^1 & 0 & B^3 & -B^2 \\ -E^2 & -B^3 & 0 & B^1 \\ -E^3 & B^2 & -B^1 & 0 \end{bmatrix} \begin{bmatrix} -1 & & & \\ & 1 & & \\ & & 1 & \\ & & & 1 \end{bmatrix}$$

This gives the homogeneous Maxwell Eq^s. Not entirely obvious... let's break it down

Consider $\mu=0, \nu=0, \lambda=i$

$$\partial_0 F_{0i} + \partial_0 F_{i0} + \partial_i F_{00} = 0$$

tautology · $F_{0i} = -F_{i0}$

Consider $\mu=0, \nu=i, \lambda=j$

$$\partial_0 F_{ij} + \partial_i F_{j0} + \partial_j F_{0i} = 0$$

$$\partial_0 \tilde{\epsilon}^{ijk} B^k + \partial_i (+E^j) + \partial_j (-E^i) = 0$$

$$F_{0i} = -E^i = -F_{i0}$$

$$i=1, j=2 \quad \partial_0 B^3 + \partial_1 E^2 - \partial_2 E^1 = 0$$

$$(\nabla \times \vec{E})^3 + \partial_0 B^3 = 0$$

$$F_{ij} = \tilde{\epsilon}^{ijk} B^k$$

Recall that $\epsilon = \delta - \mathcal{A}$, $\epsilon^{ijk} \epsilon_{kab} = \delta_{ia} \delta_{jb} - \delta_{ib} \delta_{ja}$

$$\partial_0 (\epsilon^{abi} \epsilon^{ijk} B^k) + \partial_i \epsilon^{abi} E^j - \partial_j \epsilon^{abi} E^i = 0$$

$$\partial_0 (\delta_{aj} \delta_{bk} - \delta_{ak} \delta_{bj}) B^k + \epsilon^{abi} \partial_i E^j - \epsilon^{abi} \partial_j E^i = 0$$

Anyway... there are better methods...

Energy and Momentum: from particles to stress-energy tensor

(3)

We've done some of this in LECTURE 7, but I summarize §1.9 of Carroll here:

$X^\mu(\tau)$ worldline of massive particle parametrized by proper time τ

$U^\mu = \frac{dX^\mu}{d\tau}$: 4-velocity $U^\mu U_\mu = -1$ (Carroll sets $c=1$)

$P^\mu = m U^\mu$: 4-momentum $P^\mu P_\mu = -m^2$ $\not\equiv E = \sqrt{m^2 + p^2}$

$F^\mu = m \frac{d^2 X^\mu}{d\tau^2} = \frac{dP^\mu}{d\tau}$

$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$ generalizes to $F^\mu = -q U^\lambda F_\lambda{}^\mu$ (tensorial eqⁿ)

Extended systems made of many particles can be described as a fluid which is characterized by density, pressure, entropy, viscosity etc...

$T^{\mu\nu}$ denotes the energy-momentum tensor
 it is a symmetric (2,0)-tensor.
 "the flux of 4-momentum P^μ across surface of constant X^ν "

T^{00} is flux of P^0 in the X^0 direction, the rest frame energy density

$T^{0i} = T^{i0}$ is the momentum density

$T^{ij} = T^{ji}$ gives forces between neighboring fluid elements, that is stress

- T^{ii} gives pressure
- $T^{i\hat{j}}$ ($i \neq \hat{j}$) gives shearing forces

All comments for rest frame of a fluid

DUST: Collection of particles at rest w.r.t. each other

$\mathcal{U}^\mu(x)$ should be constant as x ranges over the dust.

$N^\mu = n \mathcal{U}^\mu$: "number-flux four-vector"
(n = # density of particles as measured in their rest frame.)

Since $\mathcal{U}^\mu = (\gamma(u)c, \gamma(u)\vec{u})$ in the rest frame $\vec{u} = 0$, $\gamma = 1$ and $c = 1$ (carr.)
then $N^0 = n \mathcal{U}^0 = n$ in the rest-frame.

N^0 measured in another frame gives # density in that frame
 N^i measures flux of particles in X^i -direction

Assume all the particles have mass m then

$\rho = mn$ = rest frame energy density

But, $P^\mu = (m, 0, 0, 0)$ and $N^\mu = (n, 0, 0, 0)$ are vectors in the rest frame
for which $P^\mu \otimes N^\nu$ has $(0, 0)$ -component (mn) of energy density... so...

$$T_{\text{dust}}^{\mu\nu} = P^\mu N^\nu = mn \mathcal{U}^\mu \mathcal{U}^\nu = \rho \mathcal{U}^\mu \mathcal{U}^\nu$$

← stress-energy tensor for dust.

PERFECT FLUID

- constant rest frame energy density ρ
- isotropic rest frame pressure p

$$T^{\mu\nu} = \begin{bmatrix} \rho & 0 & 0 & 0 \\ 0 & p & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & p \end{bmatrix}$$

← Stress-Energy Tensor of perfect fluid in its rest frame

What is the general form of $T^{\mu\nu}$ though?

Since $T^{\mu\nu}_{\text{DUST}} = \rho U^\mu U^\nu$

$$T^{\mu\nu}_{\text{perfect fluid}} = (\rho + p) U^\mu U^\nu = \begin{bmatrix} \rho + p & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$$

we need to obtain the boxed eqⁿ for the right formulation ...

$$T^{\mu\nu} = (\rho + p) U^\mu U^\nu + p \eta^{\mu\nu}$$

← tensorial and reduces to $T^{\mu\nu}$ in rest frame for perfect fluid as we want.

Defⁿ relating pressure p to energy density ρ is known as an equation of state; $p = p(\rho)$

- DUST: $p = 0$
- ISOTROPIC GAS OF PHOTONS: $p = \frac{1}{3} \rho$
- VACUUM ENERGY: $p_{\text{vac}} = -\rho_{\text{vac}}$

CONSERVATION LAWS

$$\partial_\mu T^{\mu\nu} = 0$$

← conservation of stress-energy tensor so 4-dim'l divergence. We should unpack this physical equation.

Basically, $\nu = 0$ gives conservation of energy whereas $\nu = 1, 2, 3$ describes conservation of momentum.

PERFECT FLUID: $T^{\mu\nu} = (p + \rho) U^\mu U^\nu + p \eta^{\mu\nu}$

$$\partial_\rho T^{\rho\nu} = (p + \rho) [\partial_\rho U^\rho U^\nu + U^\rho \partial_\rho U^\nu] + p \partial_\rho (\underbrace{\eta^{\rho\nu}}_{= 2})$$

0 in flat spacetime η is constant matrix
 Oops, I forgot ρ, p aren't constants)

Notice $U^\rho U_\rho = -1 \Rightarrow \partial_\nu (U^\rho U_\rho) = 0$

$$\text{or } (\partial_\nu U^\rho) U_\rho + U^\rho (\partial_\nu U_\rho) = 0 \Rightarrow U_\nu \partial_\rho U^\nu = 0$$

$$U_\nu \partial_\rho T^{\rho\nu} = (p + \rho) U_\nu U^\nu \partial_\rho U^\rho + 0 \Rightarrow \partial_\rho (p + \rho) U^\rho U^\nu + (\partial_\rho p) U_\nu \eta^{\rho\nu}$$

$$= -(p + \rho) \partial_\rho U^\rho - \partial_\rho (p + \rho) U^\rho + (\partial_\rho p) U^\rho$$

$$= -\partial_\rho (p + \rho) U^\rho = 0$$

← some simplifications

Ok, now take non-relativistic limit where $U^\rho = (1, v^i)$, $|v^i| \ll 1$, $p \ll \rho$

$$\partial_t \rho + \nabla \cdot (\rho \vec{v}) = 0$$

← continuity of energy density

Remark: Carroll then looks at \perp to U^ρ -direction and derives Euler's equation from fluid mechanics.