

## BEYOND CONSTANT COEFFICIENTS; VARIABLE COEFFICIENTS ETC...: (§4.7)

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You may have noticed the theory discussed on (61) → (66) largely applies to the variable coefficient problem,

$$a(x)y''(x) + b(x)y'(x) + c(x)y(x) = g(x) \quad (*)$$

We still have  $L[y] = (aD^2 + bD + c)[y] = g$ . However, since  $Da, Db, Dc \neq 0$  we can no longer factor  $L$  as we discussed on (67) → (74). It is still the case that (\*) has a unique sol<sup>n</sup> given some initial conditions  $y(x_0) = y_0$  and  $y'(x_0) = y_1$ , however the domain of the sol<sup>n</sup> will not necessarily extend to  $\mathbb{R}$ . We can only hope for solutions to extend over domains where  $a, b, c, g$  are continuous. When a sol<sup>n</sup> of (\*) exists it still has the general form

$$y = c_1 y_1 + c_2 y_2 + y_p \quad \text{where } L[y_1] = L[y_2] = 0$$

Variation of parameters also works for (\*). The main difference is we need new methods to find  $y_1$  and  $y_2$ .

### Cauchy Euler Problems:

Let  $a, b, c$  be constants. We can solve  $at^2y'' + bty' + cy = 0$  without too much trouble. Assume the sol<sup>n</sup> has the form  $y = t^r$  then  $y' = rt^{r-1}$  and  $y'' = r(r-1)t^{r-2}$  hence we need

$$at^2 r(r-1)t^{r-2} + btrt^{r-1} + ct^r = 0$$

$$(ar(r-1) + br + c)t^r = 0, \quad \text{assume } t > 0.$$

Our guess of  $y = t^r$  will solve (CE) if the constant  $r$  satisfies the

$$\text{characteristic equation } ar(r-1) + br + c = 0$$

(text calls this the indicial eq<sup>n</sup> in §8.5)

**E87** Solve  $t^2y'' + 6ty' + 6y = 0$ . Use calculation above, we need  $r(r-1) + 6r + 6 = r^2 + 5r + 6 = (r+3)(r+2) = 0 \therefore r = -3, -2$ . The sol<sup>n</sup> has the form  $y = c_1 t^{-3} + c_2 t^{-2}$ .

**E88** Solve  $t^2 y'' + ty' + y = 0$ . Look for sol<sup>n</sup>  $y = t^r$ . The Char. Eq<sup>n</sup> reads  $r(r-1) + r + 1 = r^2 + 1 = 0 \implies r = \pm i$ .

We find sol<sup>n</sup>'s  $y = t^i$  and  $y = t^{-i}$ , notice these are complex-valued sol<sup>n</sup>'s. No problem, we have real sol<sup>n</sup>'s  $y_1 = \text{Re}(t^i)$  and  $y_2 = \text{Im}(t^i)$  hidden inside  $y = t^i$ .

Observe that:

$$t^i = e^{\ln(t^i)} = e^{i \ln(t)} = \cos(\ln(t)) + i \sin(\ln(t)).$$

Therefore, assuming  $t > 0$ ,

$$y_1 = \text{Re}(t^i) = \cos(\ln(t))$$

$$y_2 = \text{Im}(t^i) = \sin(\ln(t))$$

The general sol<sup>n</sup> is  $y = C_1 \cos(\ln(t)) + C_2 \sin(\ln(t))$

**E89** Solve  $y'' - \frac{1}{x}y' + \frac{1}{x^2}y = 0$ . Again we search for sol<sup>n</sup>'s of the form  $y = t^r$  by solving Char. Eq<sup>n</sup>  $r(r-1) - r + 1 = 0$  which yields  $(r-1)^2 = 0$  and so  $r = 1$ . Hence  $y_1 = t$

is a sol<sup>n</sup>. BUT, what is  $y_2$ ? We know the general sol<sup>n</sup> will have the form  $y = C_1 t + C_2 y_2$ .

We are unable to solve this w/o a way to find  $y_2$ !

Fortunately, reduction of order gives us a technique to find a 2<sup>nd</sup> linearly independent sol<sup>n</sup>  $y_2$  if we already have  $y_1$ . (we'll prove this formula on 96 but for now let's use it)

$$y_2 = y_1 \int \frac{e^{-\int P dt}}{y_1^2} dt = t \int \frac{e^{\int \frac{1}{x} dt}}{t^2} dt = t \int \frac{e^{\ln(x)}}{t^2} dt = t \int \frac{dt}{t}$$

Thus we find  $y_2 = t \ln(t)$  and  $y = C_1 t + C_2 t \ln(t)$   
general solution.

Remark: this is not an isolated occurrence. If we get double root of  $r$  then the Cauchy-Euler Equation has sol<sup>n</sup>  $y = C_1 t^r + C_2 t^r \ln(t)$  for  $t > 0$

## REDUCTION OF ORDER: HOW TO FIND $y_2$ FROM $y_1$ :

(96)

Suppose we have a nonzero sol<sup>n</sup>  $y_1$  of  $y'' + py' + qy = 0$  where  $p, q$  could be functions. Our goal is to find a second linearly independent sol<sup>n</sup>  $y_2$ . Since  $y_1, y_2$  need to be LI if we have  $y_2/y_1 = v$  then  $v$  cannot be constant. (otherwise  $y_1, y_2$  would be linearly dependent). Consider then,

$$y_2 = v y_1 \quad \text{where} \quad y_1'' + p y_1' + q y_1 = 0 \quad (\text{since } y_1 \text{ a sol}^n)$$

Differentiate,

$$y_2' = v' y_1 + v y_1'$$

$$y_2'' = v'' y_1 + v' y_1' + v' y_1' + v y_1'' = v'' y_1 + 2v' y_1' + v y_1''$$

We require  $y_2'' + p y_2' + q y_2 = 0$ ,

$$v'' y_1 + 2v' y_1' + v y_1'' + p(v' y_1 + v y_1') + q v y_1 = 0$$

$$v \underbrace{(y_1'' + p y_1' + q y_1)}_{\text{Zero since } y_1 \text{ is a sol}^n} + v'' y_1 + 2v' y_1' + p v y_1' = 0$$

$$y_1 v'' + (2y_1' + p y_1) v' = 0$$

Let  $w = v'$  then  $v'' = w'$  and we find  $y_1 \frac{dw}{dt} + (2y_1' + p y_1) w = 0$

$$\frac{dw}{dt} + \left( \frac{2y_1'}{y_1} + p \right) w = 0$$

We can use integrating factor technique,

$$\begin{aligned} \mu &= \exp\left(\int \left(\frac{2y_1'}{y_1} + p\right) dt\right) = \exp\left(\int \frac{2dy_1}{y_1}\right) \exp\left(\int p dt\right) \\ &= \exp(2 \ln|y_1|) \exp\left(\int p dt\right) \\ &= \exp(\ln|y_1|^2) \exp\left(\int p dt\right) \\ &= y_1^2 e^{\int p dt} \end{aligned}$$

Multiplying by integrating factor & use product rule to find that,

$$\frac{d}{dt}(\mu w) = 0 \Rightarrow y_1^2 e^{\int p dt} w = k = \text{constant}$$

## Continuing derivation of reduction of order formula:

(97)

Recall we have  $w = \frac{dv}{dt}$  where  $y_2 = v y_1$ , and  $y_1$  was given whereas  $y_2$  we wish to calculate from  $y_1$  somehow. Furthermore, we just derived that

$$y_1^2 e^{\int p dt} w = k$$

$$\Rightarrow \frac{dv}{dt} = \frac{k}{y_1^2 e^{\int p dt}}$$

$$\Rightarrow v = \int \frac{k e^{-\int p dt}}{y_1^2} dt$$

$$\Rightarrow \frac{y_2}{y_1} = \int \frac{k e^{-\int p dt}}{y_1^2} dt$$

$$\therefore \boxed{y_2 = y_1 \int \frac{\exp(-\int p dt)}{y_1^2} dt}$$

We can choose  $k = 1$  for convenience, any other nonzero value would also yield a linearly independent function from  $y_1$ .

## LOOKING BACK:

The reduction of order formula will derive the 2<sup>nd</sup> solution in the repeated root case for  $ay'' + by' + cy = 0$ ,

**E90**

Consider  $y'' - 2y' + y = 0$

$$\lambda^2 - 2\lambda + 1 = (\lambda - 1)^2 = 0 \quad \therefore \underline{y_1 = e^t}$$

In this case  $p = -2$  thus  $-\int p dt = \int 2 dt = 2t$  hence,

$$y_2 = e^t \int \frac{e^{2t}}{(e^t)^2} dt = e^t \int 1 dt = t e^t \quad \therefore \underline{y_2 = t e^t}$$

- This gives us a way of explaining where the mysterious  $t$  in the double root sol<sup>n</sup> came from. You can also derive the  $t$  from the matrix exponential applied to the complementary matrix system or Laplace Transforms. The method here is probably the most straight-forward justification.

ANOTHER APPROACH TO CAUCHY EULER PROBLEM:

$$\text{Consider } at^2 y'' + bt y' + cy = 0$$

Try the change of variables  $t = e^x$

That is  $x = \ln(t)$ . Notice

$$\frac{dy}{dt} = \frac{dx}{dt} \frac{dy}{dx} = \frac{1}{t} \frac{dy}{dx} = \frac{dy}{dt} \cdot (i)$$

$$\frac{d^2y}{dt^2} = \frac{d}{dt} \left[ \frac{1}{t} \frac{dy}{dx} \right] = \frac{-1}{t^2} \frac{dy}{dx} + \frac{1}{t} \frac{d}{dt} \left( \frac{dy}{dx} \right) = \frac{-1}{t^2} \frac{dy}{dx} + \frac{1}{t} \frac{dx}{dt} \frac{d}{dx} \left( \frac{dy}{dx} \right)$$

$$\Rightarrow \frac{d^2y}{dt^2} = \frac{-1}{t^2} \frac{dy}{dx} + \frac{1}{t^2} \frac{d^2y}{dx^2} \cdot (ii)$$

Substitute (i) & (ii) into the Cauchy Euler Problem,

$$at^2 \left( \frac{-1}{t^2} \frac{dy}{dx} + \frac{1}{t^2} \frac{d^2y}{dx^2} \right) + bt \left( \frac{1}{t} \frac{dy}{dx} \right) + cy(x) = 0$$

$$\Rightarrow a \frac{d^2y}{dx^2} + (b-a) \frac{dy}{dx} + cy(x) = 0$$

We'll find solutions of the form

$$\text{I.) } y(x) = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x} \quad \xrightarrow{x = \ln(t)} \quad y = c_1 t^{\lambda_1} + c_2 t^{\lambda_2}$$

$$\text{II.) } y(x) = c_1 e^{\lambda x} + c_2 x e^{\lambda x} \quad \xrightarrow{x = \ln(t)} \quad y = c_1 t^{\lambda} + c_2 t^{\lambda} \ln(t)$$

$$\text{III.) } y(x) = c_1 e^{\alpha x} \sin \beta x + c_2 e^{\alpha x} \cos \beta x \quad \xrightarrow{x = \ln(t)} \quad y = c_1 t^{\alpha} \sin(\beta \ln(t)) + c_2 t^{\alpha} \cos(\beta \ln(t))$$

- Why is the Cauchy-Euler Problem so similar to the constant coefficient case?
- See §8.5 for the text's treatment of the Cauchy Euler Problem,

# On LINEAR DIFFERENTIAL EQUATIONS AND EIGENFUNCTIONS

(99)

Let  $T$  be a linear operator on smooth functions on  $I \subseteq \mathbb{R}$  means  $T: C^\infty(I) \rightarrow C^\infty(I)$  and it satisfies

(1.)  $T(cf) = cT(f)$ , for all  $f \in C^\infty(I)$  and  $c \in \mathbb{R}$

(2.)  $T(f+g) = T(f) + T(g)$ , for all  $f, g \in C^\infty(I)$ .

**Eq 1** A constant coefficient ODE has the following form:

$$\underbrace{\frac{d^n}{dx^n}(y) + P_n \frac{d^{n-1}}{dx^{n-1}}(y) + \dots + P_2 \frac{d}{dx}(y) + P_1 y}_{P(D)[y]} = 0$$

$$P(D)[y] = 0$$

Where  $P(D) = D^n + P_n D^{n-1} + \dots + P_2 D + P_1$ . The nice feature of this smooth linear operator is that it factored completely into

$$P(D) = L_1 \circ L_2 \circ \dots \circ L_k$$

where  $k \leq n$  and  $L_1, L_2, \dots, L_k$  are linear operators.

In particular  $L_j = (D + \lambda_j)^{m_j}$  or  $L_j = [(D - \alpha_j)^2 + \beta_j^2]^{m_j}$  for

each  $j = 1, 2, \dots, k$ . Fortunately if  $L_j(f) = 0$  for

any particular  $j$  then  $P(D)(f) = 0$ . This

result depended strongly on the fact we can

commute the factors in  $P(D)$ ,

$$P(D) = L_1 \circ L_2 \circ \dots \circ L_j \circ \dots \circ L_k = L_1 \circ L_2 \circ \dots \circ L_k \circ L_j$$

$$\begin{aligned} \text{Hence } L_j(f) = 0 &\Rightarrow P(D)(f) = L_1 \circ L_2 \circ \dots \circ L_k \circ L_j(f) \\ &= L_1 \circ L_2 \circ \dots \circ L_k(0) = 0. \end{aligned}$$

**Def<sup>n</sup>**  $f$  is an eigenfunction with eigenvalue  $\lambda$  of the linear operator  $T$  if  $T(f) = \lambda f$ . Equivalently  $(T - \lambda)(f) = 0$ . If  $(T - \lambda)^2(f) = 0$  and  $(T - \lambda)(f) \neq 0$  then  $f$  is a generalized eigenfunction of order two.

**E9a** Let  $\lambda \in \mathbb{R}$ . The solutions to the constant coefficient ODEs are eigenfunctions or generalized eigenfunctions of  $D = d/dx$ ,

$$(D - \lambda)(e^{\lambda x}) = \lambda e^{\lambda x} - \lambda e^{\lambda x} = 0. \quad \therefore f_1(x) = e^{\lambda x} \text{ eigenfnct.}$$

$$(D - \lambda)^2(xe^{\lambda x}) = 0 \quad \therefore f_2(x) = xe^{\lambda x} \text{ generalized eigenfnct. order 2.}$$

$$(D - \lambda)^3(x^2e^{\lambda x}) = 0 \quad \therefore f_3(x) = x^2e^{\lambda x} \text{ generalized eigenfnct. order 3.}$$

When  $\lambda \in \mathbb{C}$  then we find  $e^{\lambda x}, xe^{\lambda x}, \dots$  are generalized eigenfunctions of  $D$  which are complex-valued. The  $\text{Re}(x^s e^{\lambda x})$  and  $\text{Im}(x^s e^{\lambda x})$  are not eigenfunctions of  $(D - \lambda)^{s+1}$  however they are linearly independent. Eigenfunctions with distinct eigenvalues are LI and we can show generalized eigenfunctions  $e^{\lambda x}, xe^{\lambda x}, x^2e^{\lambda x}$  etc... form a LI set of functions.

Remark: for notational simplicity assume the roots of a constant coefficient ODE  $L[y] = 0$  are all real

$\lambda_1, \lambda_2, \dots, \lambda_k$  with multiplicities  $m_1, m_2, \dots, m_k$  then

$$L[y] = \underbrace{(D - \lambda_1)^{m_1} (D - \lambda_2)^{m_2} \dots (D - \lambda_k)^{m_k}}_{\text{polynomial in } D} [y]$$

polynomial in  $D$ .

made of commuting factors

①  $D$  commutes with  $D$

②  $D$  commutes with  $\lambda$

③  $\lambda_1$  commutes with  $\lambda_2$

$$(D - \lambda_1)(D - \lambda_2)(f) = (D - \lambda_1)(f' - \lambda_2 f) = f'' - \lambda_2 f' - \lambda_1 f' + \lambda_1 \lambda_2 f$$

$$(D - \lambda_2)(D - \lambda_1)(f) = (D - \lambda_2)(f' - \lambda_1 f) = f'' - \lambda_1 f' - \lambda_2 f' + \lambda_2 \lambda_1 f$$

Clearly  $(D - \lambda_1)(D - \lambda_2)(f) = (D - \lambda_2)(D - \lambda_1)(f)$  for all  $f$  thus the operators commute.

I'll cut straight to the point here

$$ax^2y'' + bxy' + cy = 0$$

$$\underbrace{(a(xD)^2 + b(xD) + c)}_{\text{polynomial in } xD} [y] = 0$$

polynomial in  $xD$

$\Rightarrow$  can factor into smaller polynomials in  $xD$ .

$\Rightarrow$  can break down problem to finding eigenfunctions of  $xD$  instead.

Notice  $y = x^r$  is an eigenfunction of  $xD$  since  $xD(x^r) = xr x^{r-1} = r x^r$

$\Rightarrow y = x^r$  is eigenfunction of  $xD$  with eigenvalue  $r$ .

How do we factor  $(xD)^2 + b(xD) + c$ ?

It's not as simple as the constant coefficient case. Let's work backwards to see how it goes,

$$\begin{aligned} (xD + \alpha)(xD + \beta)[f] &= (xD + \alpha)[xf' + \beta f] \\ &= xD(xf') + xD(\beta f) + \alpha xf' + \alpha \beta f \\ &= xf' + x^2 f'' + \beta xf' + \alpha xf' + \alpha \beta f \\ &= (x^2 D^2 + (\alpha + \beta + 1)x D + \alpha \beta)[f] \end{aligned}$$

**E93** Factor  $L = x^2 D^2 + 4xD + 2$ . I identify  $\alpha = 1, \beta = 2$

thus  $L = (xD + 1)(xD + 2)$ . We can solve

$$x^2 y'' + 4xy' + 2y = 0 \Rightarrow (xD + 1)(xD + 2)[y] = 0$$

$\downarrow$  eigenfunction  $y_1 = x^{-1}$        $\downarrow$  eigenfunction  $y_2 = x^{-2}$

$$\therefore y = C_1/x + C_2/x^2$$

**Remark:** Don't worry I'm majorly digressing for the last 4 pages.