

LECTURE 7

- we study the kinematics of circular motion. Tangential and normal accelerations are defined and discussed. To begin we provide a little math background.

Thm / Let \vec{A}, \vec{B} be vector-valued functions of time with differentiable component functions then

$$\frac{d}{dt} (\vec{A} \cdot \vec{B}) = \frac{d\vec{A}}{dt} \cdot \vec{B} + \vec{A} \cdot \frac{d\vec{B}}{dt}$$

Proof:

$$\begin{aligned} \frac{d}{dt} (\vec{A} \cdot \vec{B}) &= \frac{d}{dt} \left[\sum_{j=1}^3 A_j B_j \right] && : \text{def}^n \text{ of dot-product} \\ &= \sum_{j=1}^3 \frac{d}{dt} (A_j B_j) && : \text{linearity of } \frac{d}{dt} \\ &= \sum_{j=1}^3 \left[\frac{dA_j}{dt} B_j + A_j \frac{dB_j}{dt} \right] && : \text{product rule on each } j. \\ &= \sum_{j=1}^3 \frac{dA_j}{dt} B_j + \sum_{j=1}^3 A_j \frac{dB_j}{dt} \\ &= \frac{d\vec{A}}{dt} \cdot \vec{B} + \vec{A} \cdot \frac{d\vec{B}}{dt} // \end{aligned}$$

Thm / Let $\vec{A}, \vec{B}, \vec{C}$ be vectors and $\alpha, \beta \in \mathbb{R}$ then

(1) $\vec{A} \cdot \vec{B} = \vec{B} \cdot \vec{A}$, (2) $\vec{A} \cdot (\vec{B} + \vec{C}) = \vec{A} \cdot \vec{B} + \vec{A} \cdot \vec{C}$, $\vec{A} \cdot \alpha \vec{B} = \alpha \vec{A} \cdot \vec{B}$

Proof: Einstein's sneaky notation $\vec{A} \cdot \vec{B} = A_i B_i$ omits $\sum_{i=1}^3$.
I'll demonstrate it here for fun:

- (1) $\vec{A} \cdot \vec{B} = A_i B_i = B_i A_i = \vec{B} \cdot \vec{A}$.
- (2) $\vec{A} \cdot (\vec{B} + \vec{C}) = A_i (\vec{B} + \vec{C})_i = A_i (B_i + C_i) = A_i B_i + A_i C_i = \vec{A} \cdot \vec{B} + \vec{A} \cdot \vec{C}$.
- (3) $\vec{A} \cdot (\alpha \vec{B}) = A_i (\alpha B)_i = A_i \alpha B_i = \alpha A_i B_i = \alpha \vec{A} \cdot \vec{B}$.

Einstein's subtle notation allows focus on the essential points of the calculation. The brevity & clarity is especially important to theoretical physics calculations... //

CIRCULAR MOTION

If $t \mapsto \vec{r}(t)$ is a path around a circle with center \vec{c} and radius R we have:

$$\|\vec{r}(t) - \vec{c}\| = R \quad \forall t \in \text{dom}(\vec{r})$$

However, recall $\|\vec{A}\|^2 = \vec{A} \cdot \vec{A}$ thus,

$$(\vec{r} - \vec{c}) \cdot (\vec{r} - \vec{c}) = R^2$$

$$\Rightarrow \vec{r} \cdot \vec{r} - \vec{r} \cdot \vec{c} - \vec{c} \cdot \vec{r} + \vec{c} \cdot \vec{c} = R^2$$

$$\Rightarrow \vec{r} \cdot \vec{r} - 2\vec{r} \cdot \vec{c} + \vec{c} \cdot \vec{c} = R^2$$

Let's differentiate, note R^2, \vec{c} are constant thus \Rightarrow

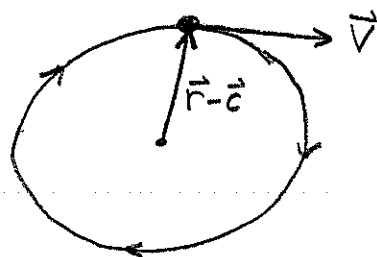
$$0 = \frac{d\vec{r}}{dt} \cdot \vec{r} + \vec{r} \cdot \frac{d\vec{r}}{dt} - 2\frac{d\vec{r}}{dt} \cdot \vec{c} \quad \left(\frac{d\vec{c}}{dt} = 0, \frac{d(R^2)}{dt} = 0 \right)$$

$$\Rightarrow 0 = 2\frac{d\vec{r}}{dt} \cdot \vec{r} - 2\frac{d\vec{r}}{dt} \cdot \vec{c}$$

$$\Rightarrow 0 = \frac{d\vec{r}}{dt} \cdot (\vec{r} - \vec{c})$$

$$\Rightarrow 0 = \vec{v} \cdot (\vec{r} - \vec{c})$$

velocity \uparrow radial vector



Differentiate once more,

$$0 = \frac{d\vec{v}}{dt} \cdot (\vec{r} - \vec{c}) + \vec{v} \cdot \left(\frac{d\vec{r}}{dt} \right)$$

$$\Rightarrow \boxed{\vec{a} \cdot (\vec{r} - \vec{c}) = -\vec{v} \cdot \vec{v}} \quad (*)$$

Let $\vec{r} = \vec{r} - \vec{c}$ and $\hat{r} = \frac{1}{r} \vec{r}$ as usual then (*) reads $\vec{a} \cdot \vec{r} = -\vec{v} \cdot \vec{v}$. Divide by r to find,

$$\vec{a} \cdot \hat{r} = \frac{-\vec{v} \cdot \vec{v}}{r} = \frac{-v^2}{r}$$

$$\therefore \boxed{a_r = \frac{-v^2}{r}}$$

\leftarrow the center seeking or centripetal acceleration must be proportional to $\frac{v^2}{r}$ for circular motion.

CIRCULAR MOTION: TAKE TWO

Let $\vec{r}(t) = R \cos(\theta) \hat{i} + R \sin(\theta) \hat{j}$. We assume circle centered at origin and we choose coordinates such that $\vec{r}(0) \propto \hat{i}$. We calculate,

θ generally function of time.

$$\vec{v} = \frac{d\vec{r}}{dt} = -R \sin\theta \frac{d\theta}{dt} \hat{i} + R \cos\theta \frac{d\theta}{dt} \hat{j}$$

$$\vec{a} = \frac{d\vec{v}}{dt} = -R \cos\theta \left(\frac{d\theta}{dt}\right)^2 \hat{i} - R \sin\theta \frac{d^2\theta}{dt^2} \hat{i} - R \sin\theta \left(\frac{d\theta}{dt}\right)^2 \hat{j} + R \cos\theta \frac{d^2\theta}{dt^2} \hat{j}$$

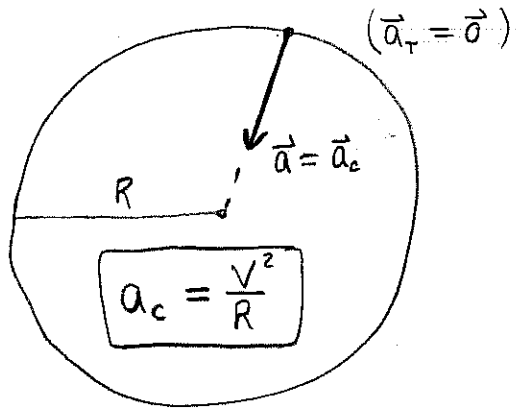
$$\Rightarrow \vec{a} = R \left(\frac{d\theta}{dt}\right)^2 \underbrace{[-\cos\theta \hat{i} - \sin\theta \hat{j}]}_{\vec{N} \text{ Unit-normal}} + R \frac{d^2\theta}{dt^2} \underbrace{[-\sin\theta \hat{i} + \cos\theta \hat{j}]}_{\vec{T} \text{ unit-tangent}}$$

Notice $s = R\theta$ for circle and so $\frac{ds}{dt} = R \frac{d\theta}{dt}$ which tells us that $v = R \frac{d\theta}{dt}$ hence,

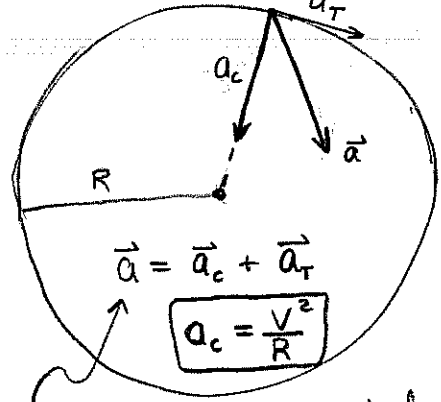
$$\vec{a} = \underbrace{\frac{v^2}{R} \vec{N}}_{\text{centripetal center-seeking acceleration}} + \underbrace{R \frac{d^2\theta}{dt^2} \vec{T}}_{\text{tangential acceleration, only nontrivial for } \frac{dv}{dt} = R \frac{d^2\theta}{dt^2} \neq 0.}$$

E1

Constant speed, $\vec{a} = \vec{a}_c \vec{N}$



Nonconstant Speed



resultant of centripetal and normal vector components.

E2 Find the centripetal acceleration at the equator of earth due to the earth's rotation. I idealize earth as sphere which spins about axis.

I take a long walk and learn $R_{\text{earth}} = 6371 \text{ km}$.

We all know that the earth rotates once per day.

It follows the speed of the comoving frame at the equator is

$$v = \frac{2\pi R_{\text{EARTH}}}{(24)(3600)\text{s}} = 0.4633 \frac{\text{km}}{\text{s}} = \underbrace{463.3 \frac{\text{m}}{\text{s}}}_{1036.6 \text{ mph}}$$

Thus you might expect a large a_c

BUT R_{EARTH} is large thus,

$$a_c = \frac{v^2}{R_E} = \frac{(463.3 \frac{\text{m}}{\text{s}})^2}{6371 \text{ km}} = \underbrace{0.034 \text{ m/s}^2}$$

$$a_c \approx 0.0034 g$$

↑
quite small in comparison to gravitational acceleration.

E3 a_c for earth relative rotation around the sun is even smaller

$$v_{\text{orb.}} = \frac{2\pi R_{\text{ORBIT}}}{(365)(86400\text{s})} = \frac{(2\pi)(1.5 \times 10^{11} \text{ m})}{(365)(86400)} = 29885.8 \frac{\text{m}}{\text{s}}$$

$$a_c = \frac{(29885.8 \text{ m/s})^2}{1.5 \times 10^{11} \text{ m}} = 0.00595 \text{ m/s}^2.$$

